

Séminaire MOCO (IRMA) 23/03/21

ϕ -FEM :

A fictitious domain method for finite element methods
on domains defined by level-sets

M. Duprez¹, V. Lleras² and A. Lozinski³,

¹Inria, Strasbourg, France

²IMAG, Montpellier, France

³Laboratoire de Mathématiques de Besançon, France



- The "standard" Finite Element Method
- Geometrical constraint on the mesh
- Previous Fictitious Domain Methods
- ϕ -FEM
- Conclusions and perspectives

Strong formulation of the Poisson equation

$$\left\{ \begin{array}{l} \text{Find } u \in H^2(\Omega) \text{ s.t. :} \\ -\Delta u = f, \text{ in } \Omega \\ u = 0, \text{ on } \partial\Omega \end{array} \right.$$



Multiplying by a "test" function v & by integration by part, one has

Weak formulation

$$\left\{ \begin{array}{l} \text{Find } u \in H_0^1(\Omega) \text{ s.t. :} \\ \Leftrightarrow \int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} f v \quad \forall v \in H_0^1(\Omega). \end{array} \right.$$

Standard FEM formulation

$$\left\{ \begin{array}{l} \text{Find } u_h \in V_h \text{ s.t. :} \\ \Leftrightarrow \int_{\Omega} \nabla u_h \cdot \nabla v_h = \int_{\Omega} f v_h \quad \forall v_h \in V_h. \end{array} \right.$$

where V_h is a subspace of $H_0^1(\Omega)$ of **finite dimensional**.

Finite element space

Let $V_h = \langle \psi_k \in H_0^1(\Omega) : k \in \{1, \dots, N\} \rangle \subset H_0^1(\Omega)$.

Equivalence with a matrix system

$$\left\{ \begin{array}{l} \text{Find } u_h = \sum U_{hk} \psi_k \in V_h \text{ s.t. :} \\ \int_{\Omega} \nabla u_h \cdot \nabla \psi_k = \int_{\Omega} f \psi_k \quad \forall k \end{array} \right. \Leftrightarrow \left\{ \begin{array}{l} \text{Find } U_h \in \mathbb{R}^N \text{ s.t. :} \\ A_h U_h = F_h \end{array} \right.$$

where

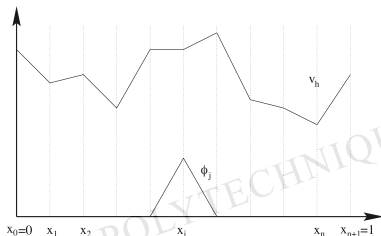
$$\left\{ \begin{array}{l} A_h = (\int_{\Omega} \nabla \psi_k \cdot \nabla \psi_j)_{kj} \\ F_h = (\int_{\Omega} f \psi_k)_k \\ U_h = (U_{hk})_k \end{array} \right.$$

Final solution

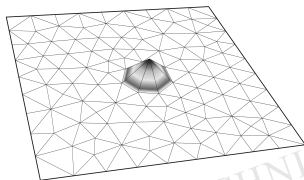
$$u_h = \sum U_{hk} \psi_k.$$

Lagrange continuous finite element

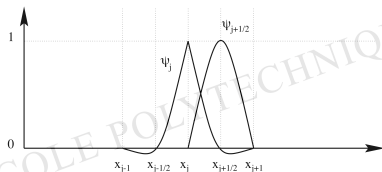
FE space : $V_h = \{\text{cont. piecewise pol. functions on a **regular mesh**\}$



Piecewise linear Lagrange FE order 1 (\mathbb{P}_1)



Order 1 (\mathbb{P}_1), dimension 2



Order 2 (\mathbb{P}_2), dimension 1

Why use polynomials ?

because the computation of the coefficient $\int_{\Omega} \nabla \psi_k \cdot \nabla \psi_l$ of the FE matrix is explicit and **exact**

Important : if ψ_k is not polynomial, this computation is not exact anymore

Why a mesh ?

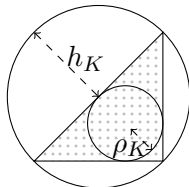
because we can **not good approximate a function** by a polynomial on the whole domain

Regular mesh : Geometrical condition

- **Simplex (triangle, tetrahedron)**

The standard FEM works under the **Ciarlet condition** [Ciarlet 78']

$$\frac{h_K}{\rho_K} < \gamma$$



- **In general (hexahedron,...)**

The standard FEM works if the Jacobian matrix of the transformation to an reference element is definite positive.

Important remark : If the mesh contains degenerated cells :

- It is not guaranty that standard FEM converges
- The **conditioning number** of the FE matrix is bad

What is the conditioning number ?

Conditioning number of a matrix A :

$$\text{Cond}(\mathbf{A}) := \|\mathbf{A}\|_2 \|\mathbf{A}^{-1}\|_2$$

For symmetric definite positive matrices :

$$\text{Cond}(\mathbf{A}) := \frac{\text{biggest eigen value}}{\text{smallest eigen value}} > 1$$

Example

$$\text{if } \mathbf{A} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \varepsilon \end{pmatrix} \text{ then } \text{Cond}(\mathbf{A}) = 1/\varepsilon.$$

If ε goes to zero, \mathbf{A} goes to a non-invertible matrix

Conclusion

the conditioning number measures the **invertibility of the matrix**.

Poor conditioning of the system matrix

Proposition

- **Non-degenerated mesh** : Suppose that the mesh satisfies the geometrical conditions, then

$$\text{Cond}(\mathbf{A}) \leq C/h^2$$

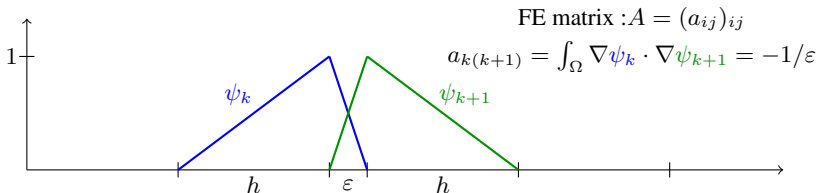
where h is the size of the cells.

- **Only one degenerated cell** :

$$\text{Cond}(\mathbf{A}) \geq C/h\varepsilon$$

where ε is the size the degenerated cell.

Idea :

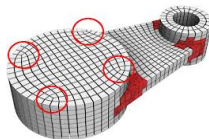


What can we do on **complex geometries** ?



In particular how to use **hexahedral meshes**

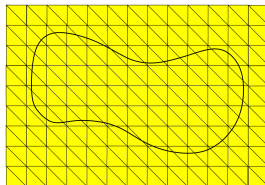
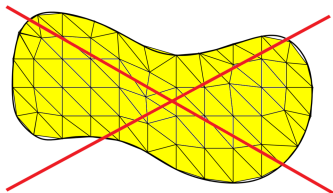
➤ Quasi-compressible elastic models : locking effect



What is the locking effect ?

Tetrahedrons have not enough degree of freedom

Alternative : Fictitious domain methods

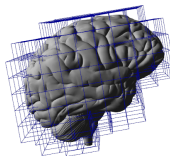


Advantages

- No need to mesh
- Regular cells

Difficulties

- Adapt the **weak formulation**
- **Conditioning** of the FE matrix



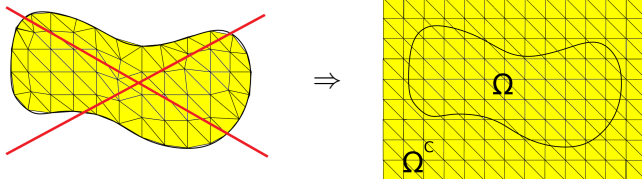
- The "standard" Finite Element Method
- Geometrical constraint on the mesh
- **Previous Fictitious domain Methods**
- ϕ -FEM
- Conclusions and perspectives

References

Saul'ev 63' (Dirichlet), Astrakhantsev 78' (Neumann),
Glowinski 92' (first proof)

Formal idea

Total energy = (energy on Ω) + $\varepsilon \times$ (energy on Ω^c) $(\varepsilon \ll 1)$



☺ **Non-conform mesh (complex and time varying geometry)**

☹ **Discontinuity, large FE matrix and bad cond. number**

☹ **Non-optimal convergence $O(\sqrt{h})$**

References

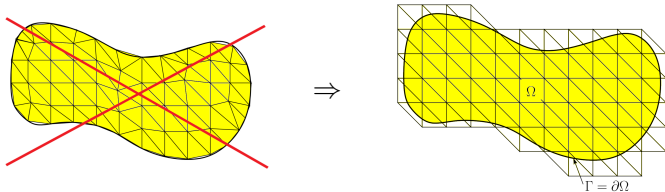
Moes-Bechet-Tourbier 2006, Haslinger-Renard 2009

Formal idea

Cut shape function : $\psi_k \longrightarrow \psi_k \mathbb{1}_\Omega$

Boundary condition on Γ : Lagrange multiplier

Conditioning of the matrix : stabilization on the boundary



☺ **Small FE matrix**

☺ **Good conditioning number**

☹ **Non-classical shape functions** and **discontinuity** in the integrals

References

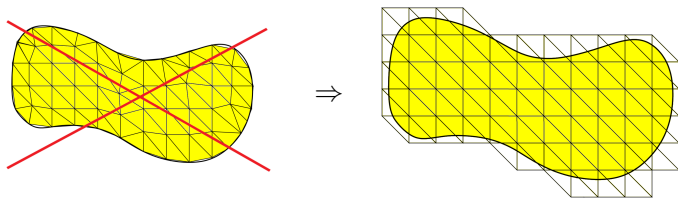
Burman et al 2010-2014

Formal idea

Partial integration on the cell near the boundary

Lagrange multiplier or penalization for the boundary conditions

Conditioning of the matrix : stabilization on the boundary



☺ **Standard shape functions**

☹ **Integral on the real boundary, cut integral.**

Formal idea (Dirichlet)

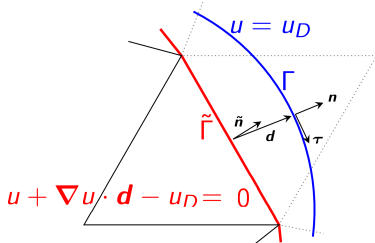
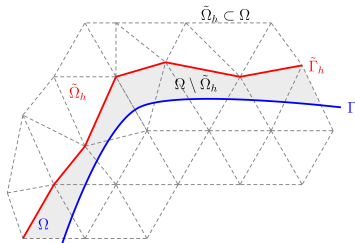
Taylor expansion of the bound. cond.

Real boundary condition on Γ : $u = u_D$

Discrete bound. cond. on $\tilde{\Gamma}$: $u + \nabla u \cdot d = u_D$

Reference

Main-Scovazzi 2017



☺ **No cut integral**

☹ **Construction of d , \mathbb{P}_k , Neumann conditions**

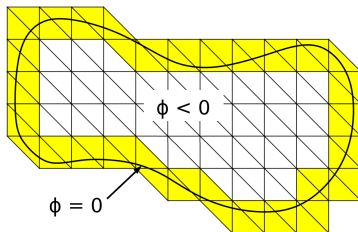
- The "standard" Finite Element Method
- Geometrical constraint on the mesh
- Previous Fictitious domain Methods
- ϕ -FEM
- Conclusions and perspectives

What is the idea of ϕ -FEM ?

Hypothesis :

Assume that Ω and Γ are given by a **level-set** function ϕ :

$$\Omega := \{\phi < 0\} \text{ and } \Gamma := \{\phi = 0\}.$$

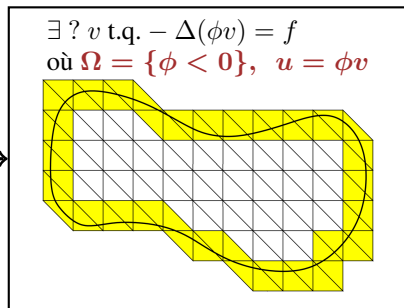
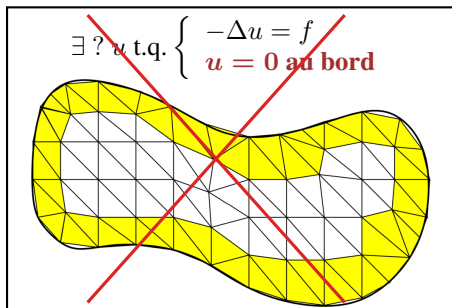


Example of level-set : signed distance

Idea of ϕ -FEM :

Include the Level-set function in the formulation
to take into account the boundary conditions

ϕ FEM for the Poisson Dirichlet Problem



ϕ -FEM : weak formulation

Hypothesis : Ω and Γ are given by a level-set function ϕ :

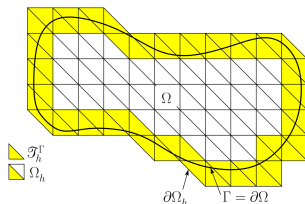
$$\Omega := \{\phi < 0\} \text{ and } \Gamma := \{\phi = 0\}.$$

ϕ -FEM formulation : Find v_h s.t.

$$\int_{\Omega_h} \nabla(\phi v_h) \cdot \nabla(\phi w_h) - \int_{\partial\Omega_h} \frac{\partial}{\partial n}(\phi v_h) \phi w_h + \text{Stab. Term} = \int_{\Omega_h} f \phi w_h \quad \forall w_h \in W_h,$$

where $W_h = \{\text{cont. piecewise pol. functions on a the mesh}\}$.

Then $u_h := \phi v_h$ as approximation of the initial problem.

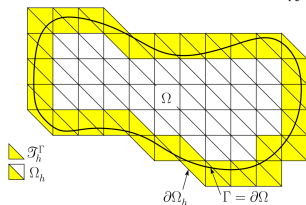


Stabilization terms (to ensure a good conditioning number)

$$\sigma h \sum_{E \in \mathcal{F}_\Gamma} \int_E \left[\frac{\partial}{\partial n}(\phi_h v_h) \right] \left[\frac{\partial}{\partial n}(\phi_h w_h) \right] + \sigma h^2 \sum_{T \in \mathcal{T}_h^\Gamma} \int_T (\Delta(\phi_h v_h) + f) \Delta(\phi_h w_h)$$

where $[\cdot]$ is the jump on the interface E

$$\begin{cases} \mathcal{T}_h^\Gamma = \{T \in \mathcal{T}_h : T \cap \Gamma_h \neq \emptyset\} \quad (\Gamma_h = \{\phi_h = 0\}); \\ \mathcal{F}_\Gamma = \{E \text{ (an internal edge of } \mathcal{T}_h) \text{ such that } \exists T \in \mathcal{T}_h : \\ T \cap \Gamma_h \neq \emptyset \text{ and } E \in \partial T\}. \end{cases}$$



Continuous problem

$$a(u, w) = l(w) \quad \forall w \in V$$

Finite element formulation

$$a_h(v_h, w_h) = l_h(w_h) \quad \forall w_h \in V_h^{(k)} \subset V$$

Theorem (D.-Lozinski 2020)

Suppose that the mesh \mathcal{T}_h is quasi-uniform (with some weak assumptions), $\Omega \subset \Omega_h$, and $f \in H^k(\Omega_h \cup \Omega)$. Let $u \in H^{k+2}(\Omega)$ be the continuous solution and $v_h \in V_h^{(k)}$ be the discret solution. Denoting $u_h := \phi_h v_h$, it holds

$$\|u - u_h\|_{1,\Omega} \leq Ch^k \|f\|_{k,\Omega_h}$$

$$\|u - u_h\|_{0,\Omega} \leq Ch^{k+1/2} \|f\|_{k,\Omega_h}.$$

with $C = C(\phi)$ and k the degree of approximation.

Remark :

Optimal order numerically

Highlights of the proof

- Extend $u \in H^{k+2}(\Omega)$ to $\tilde{u} \in H^{k+2}(\Omega_h)$ with $\tilde{u} = u$ on Ω
- **Hardy inequality**

$$\tilde{u} \in H^{k+2}(\mathcal{O}), \tilde{u}|_{\Gamma} = 0 \implies \|\tilde{u}/\phi\|_{k+1, \mathcal{O}} \leq C \|\tilde{u}\|_{k+2, \mathcal{O}}.$$

permits to properly introduce $v := (\tilde{u}/\phi) \in H^{k+1}(\Omega_h)$.

- This v satisfies

$$\int_{\Omega_h} \nabla(\phi v) \cdot \nabla(\phi w) - \int_{\partial\Omega_h} \frac{\partial}{\partial n}(\phi v) \phi w = \int_{\Omega_h} \tilde{f} \phi w, \quad \forall w \in H^1(\Omega_h)$$

with $\tilde{f} := -\Delta(\phi v) = -\Delta\tilde{u}$.

- The ghost penalty G_h permits to retrieve the coercivity of the bilinear scheme (on the discrete level) so that, introducing $e_h = v_h - I_h v$,

$$\|e_h\|_h^2 := |\phi_h e_h|_{1, \Omega_h}^2 + G_h(v_h, v_h) \lesssim |\tilde{u} - \phi_h I_h v|_{1, \Omega_h} \|e_h\|_h + \int_{\Omega_h \setminus \Omega} (f - \tilde{f}) \phi_h e_h + \dots$$

leading to an $O(h^k)$ bound since $\|f - \tilde{f}\|_{0, \Omega_h \setminus \Omega} = \|f + \Delta\tilde{u}\|_{0, \Omega_h \setminus \Omega} \lesssim h^{k-1}$,
and $\|\phi_h e_h\|_{0, \Omega_h \setminus \Omega} \lesssim h$

ϕ -FEM : Simulation of the Poisson-Dirichlet problem

Domain Ω : circle of radius $\sqrt{2}/4$ centered at $(0.5, 0.5)$

Surrounding domain : $\mathcal{O} = (0, 1)^2$

Level-set function : $\phi(x, y) = (x - 1/2)^2 - (y - 1/2)^2 - 1/8$

Exact solution : $u(x, y) = \exp(x) \times \sin(2\pi y)$

Artificial external force : $f := -\Delta u$

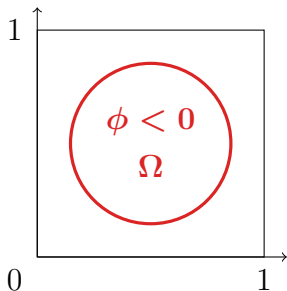
Boundary : $u_D := u(1 + \phi)$

Remark (non-homogeneous case)

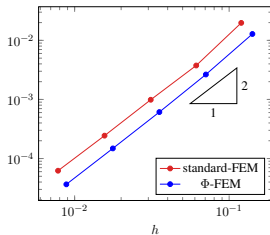
If $u := u_D$ on Γ ,

we replace $\phi_h v_h$ in the scheme

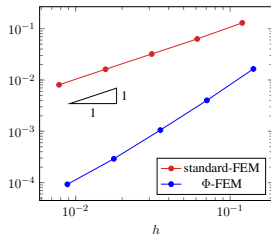
by $\phi_h v_h + u_D$ on Ω_h^Γ .



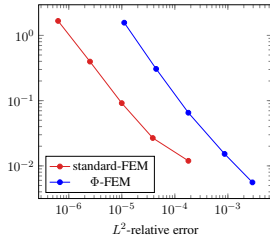
ϕ -FEM : Simulations in \mathbb{P}_1



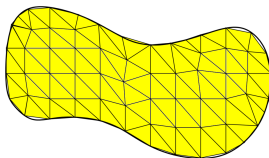
L^2 relative error



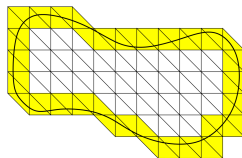
H^1 relative error



Comp. time (/error)



standard FEM

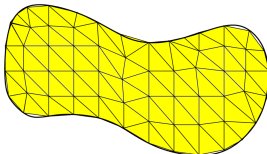


ϕ -FEM

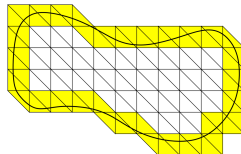
Dirichlet : $u_D = (1 + \phi)u$

Explanation of the numerical results

- Why better than standard FEM ?
 - Standard FEM : polygonal approximation of the boundary
 - ϕ -FEM : better approx. of the bound. with a levelset
- Numerical cost of ϕ -FEM :
 - Good point : small size of the FEM matrix
 - Bad point : quadrature more expensive
 - integral of polynomial product



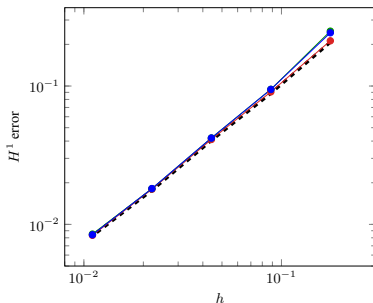
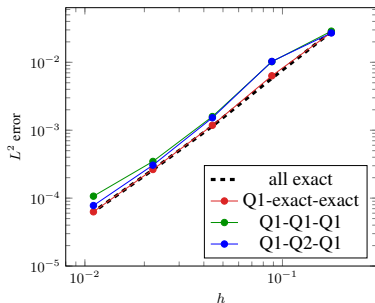
standard FEM



ϕ -FEM

Influence of an inexact quadrature

- Up to now, all the integrals were calculated exactly in all the terms of the ϕ -FEM scheme
- What if an inexact numerical quadrature is employed ?



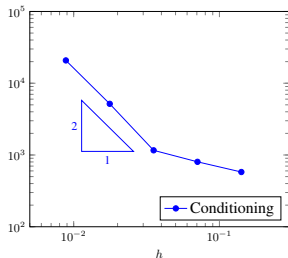
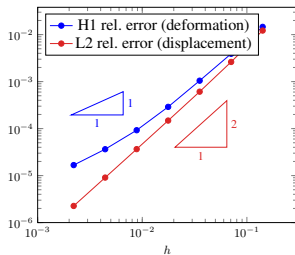
(quadrature on triangles)–(quadrature on boundary edges)–(quadrature on "ghost" edges)

Q1 : midpoint rule on a triangle or on an edge

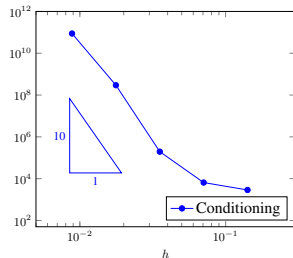
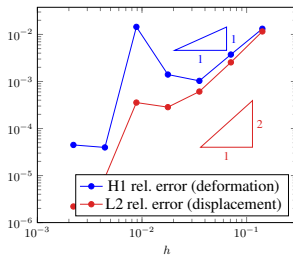
Q2 : 2 point Gaussian quadrature on an edge

ϕ -FEM : Stabilization and conditioning

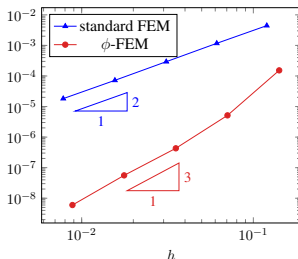
ϕ -FEM in \mathbb{P}_1 with stabilization



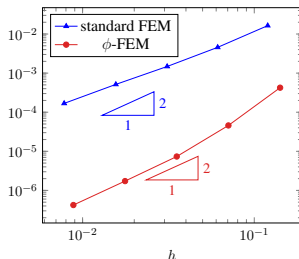
ϕ -FEM in \mathbb{P}_1 without stabilization



Remarque : ϕ -FEM works high polynomial orders



L^2 -relative error in \mathbb{P}_2



H^1 -relative error \mathbb{P}_2

Neumann boundary condition

Consider the **Poisson Neumann** problem :

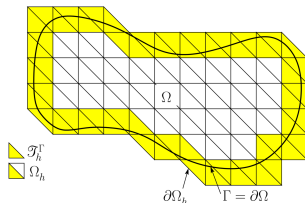
$$\begin{cases} -\Delta u + u = f, & \text{in } \Omega \\ \nabla u \cdot n = g, & \text{on } \partial\Omega \end{cases}$$

If

$$\boxed{\phi = 0 \quad \text{and} \quad \nabla\phi = n \quad \text{on } \Gamma,}$$

the last system is (formally) equivalent to

$$\begin{cases} y = -\nabla u, & \text{in } \mathcal{T}_h^\Gamma \\ \mathbf{y} \cdot \nabla\phi = p\phi + g, & \text{in } \mathcal{T}_h^\Gamma \end{cases} \approx \nabla u \cdot n = g \quad \text{on } \Gamma$$



Neumann boundary condition : weak formulation

Find $(u_h, y_h, p_h) \in W_h^{(k)}$ such that for all $(v_h, z_h, q_h) \in W_h^{(k)}$

$$\begin{aligned} \int_{\Omega_h} \nabla u_h \cdot \nabla v_h + \int_{\Omega_h} u_h v_h + \int_{\partial\Omega_h} y_h \cdot n v_h \\ + \gamma_u \int_{\Omega_h^\Gamma} (y_h + \nabla u_h) \cdot (z_h + \nabla v_h) \\ + \frac{\gamma_p}{h^2} \int_{\Omega_h^\Gamma} (y \cdot \nabla \phi_h + \frac{1}{h} p_h \phi_h - g)(z \cdot \nabla \phi_h + \frac{1}{h} q_h \phi_h) \\ + \sigma h \int_{\Gamma_h^i} [\partial_n u] [\partial_n v_h] + \gamma_{div} \int_{\Omega_h^\Gamma} (\operatorname{div} y_h + u_h - f)(\operatorname{div} z_h + v_h) \\ = \int_{\Omega_h} f v_h, \end{aligned}$$

where ϕ_h is the Lagrange interpolation of ϕ of order l and

$$\begin{aligned} W_h = \{ (u_h, y_h, p_h) \in C^0(\Omega_h) \times C^0(\Omega_h^\Gamma)^d \times L^2(\Omega_h^\Gamma) : \\ (u_h, y_h, p_h)_K \in \mathbb{P}_k \times (\mathbb{P}_k)^d \times \mathbb{P}_{k-1} \} \end{aligned}$$

Theorem (D.-Lozinski-Lleras 2021)

Suppose that the mesh is quasi-uniform (under some weak assumption), $l \geq k + 1$, $\Omega \subset \Omega_h$ and $f \in H^k(\Omega_h)$. Let $u \in H^{k+2}(\Omega)$ be the **continuous solution** and $(u_h, y_h, p_h) \in W_h^{(k)}$ be the **discrete solution**. Provided γ_{div} , γ_u , γ_p , σ are sufficiently big, it holds

$$|u - u_h|_{1,\Omega} \leq Ch^k (\|f\|_{k,\Omega_h} + \|g\|_{k+1,\Omega_h^\Gamma})$$

$$\|u - u_h\|_{0,\Omega} \leq Ch^{k+1/2} (\|f\|_{k,\Omega_h} + \|g\|_{k+1,\Omega_h^\Gamma})$$

with $C = C(\phi) > 0$.

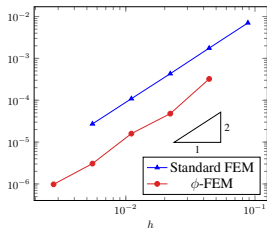
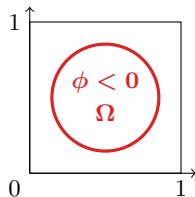
Proposition (Optimal conditioning)

The finite element matrix of ϕ -FEM satisfies

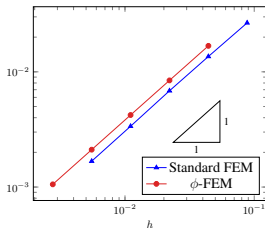
$$\text{Cond}(\mathbf{A}) \leq Ch^{-2}.$$

ϕ -FEM Neumann Poisson

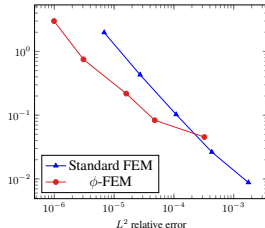
- **Level set function** :
distance to the boundary
- **Exact solution** : $u(x, y) := \sin(x) \exp(y)$
- **Source term** : $f := -\Delta u + u$
- Extrapolated **Neumann** boundary condition :
$$\tilde{g} = \frac{\nabla u \cdot \nabla \phi}{|\nabla \phi|} + u\phi$$



L^2 relative error



H^1 relative error

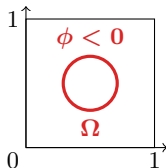


Comp. time (/error)

Stokes equations

Model : Stokes equations

$$\begin{cases} -\Delta u + \nabla p = f, & \text{in } \Omega, \\ \operatorname{div} u = 0, & \text{in } \Omega \\ u = u_D, & \text{on } \Gamma \end{cases}$$



Standard-FEM model

Find $u_h \in V_h, p_h \in M_h$ s.t. $\forall v_h \in V_h, q_h \in M_h$:

$$\int_{\Omega_h} 2Du_h : Dv_h - \int_{\Omega_h} p_h \operatorname{div} v_h - \int_{\Omega_h} q_h \operatorname{div} u_h = \int_{\Omega_h} f v_h$$

ϕ -FEM model

Find $w_h \in V_h, p_h \in M_h$ s.t. $\forall s_h \in V_h, q_h \in M_h$:

$$\begin{aligned} & \int_{\Omega_h} 2D(u_D + \phi_h w_h) : D(\phi_h s_h) - \int_{\Gamma_h} (2D(u_D + \phi_h w_h) - pI)n \cdot \phi_h s \\ & - \int_{\Omega_h} p \operatorname{div}(\phi_h s_h) - \int_{\Omega_h} q \operatorname{div}(u_D + \phi_h w_h) + \text{Stab. Term} = \int_{\Omega_h} f \phi_h s \end{aligned}$$

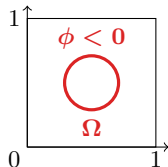
Stokes equations

Exact solution

$$u_1 = \cos(\pi x) \sin(\pi y)$$

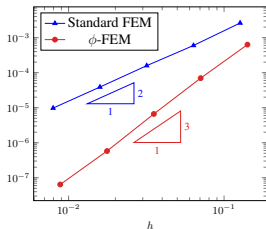
$$u_2 = \sin(\pi x) \cos(\pi y)$$

$$p = (y - 0.5) \cos(2\pi x) + (x - 0.5) \sin(2\pi y)$$

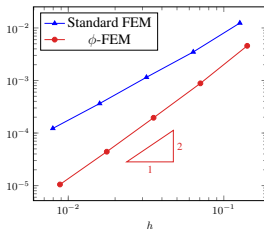


Finite element space

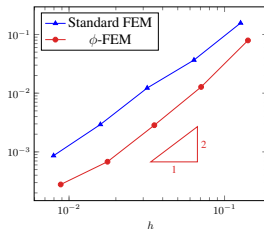
Taylor-Hood (P2-P1)



L^2 -relative error of u



H^1 -relative error of u



L^2 -relative error of p

Results

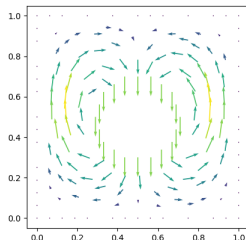
- Optimal convergence
- Discrete problem well conditioned
- Simple implementation : standard shape functions
- Formulation available for any order of approximation
- Precised and fast

MIMESIS Inria Team

Adapted to convolutional neural network

Work in progress

- Dynamic system
- Fluid structure : moving solid
- Non-linear model
- Optimization shape



Thanks for your attention !